

# **AERODYNAMIC ANALYSIS OF COMPLIANT THRUST FOIL BEARINGS**

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Mechanical  
Engineering

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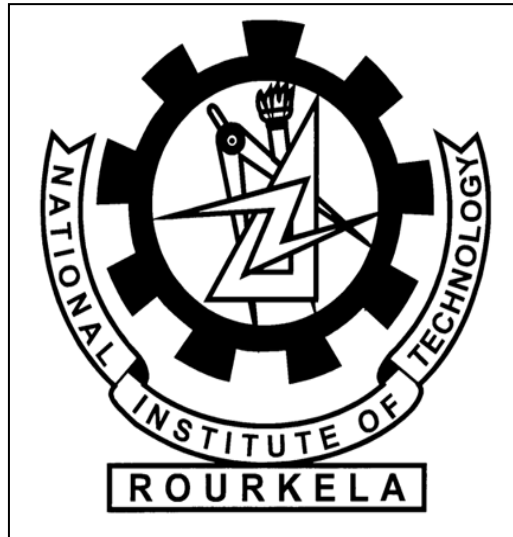
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2012-13

## **CERTIFICATE**



### **NATIONAL INSTITUTE OF TECHNOLOGY, ROURKELA**

This is to certify that this thesis titled by “**Aerodynamic Analysis of Compliant Thrust Foil Bearings**” submitted by **D. Jaswant Kumar (Roll No. : 109ME0378)** in the partial fulfilment of the requirements for the degree of **Bachelor of Technology in Mechanical Engineering, National Institute of Technology, Rourkela**, is an original and authentic work carried out by him under my supervision. To the best of my knowledge the data or matter used in this thesis has not been submitted to any other University/ Institute for the award of any degree or diploma.

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I also feel privileged to fulfil our parents' ambition and I will always be grateful for their support.

Date:

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## **ABSTRACT**

In turbo-expanders, the shaft rotates at a very high speed (i.e. 100000 RPM). At this high speeds, normal bearings cannot be used as it results in high friction and wear of the bearings. That is the reason; we are going to use gas foil bearings as a recent and advanced alternative. Here, in this project we are going to design bearings required to support the shaft which runs at a very high speed. Current project concentrate to analyse the load bearing capacity of the thrust bearings. In the analysis Reynolds' Equation is used to know the pressure distribution of these bearings. The Reynolds' Equation is solved by using FINITE DIFFERENCE METHOD and using many assumptions to know the pressure distribution of the thrust bearings. Finite Difference Method is a numerical technique by the principle of discretization to find the approximate solutions of engineering problems. The result comes after a many number of iterations based on a convergence condition. We are using MATLAB (Matrix Laboratory) software to implement Finite Difference Method to solve Reynolds' Equation. A MATLAB program is written which contains multiple loops that solves the Reynolds' Equation and gives Pressure plots. After the pressure distribution is known, load carrying capacity of the bearing is calculated and their variations with different parametes are presented. The results of this foil thrust bearings are compared with the load carrying capacity of rigid bearings. The analysis was also done for different types for foil bearings by taking different materials of the bump foil.

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## NOMENCLATURE

$A$  = Area ,  $\text{mm}^2$

$b$  = Extent of bump foil

$h$  = Film Thickness (mm)

$\bar{h}$  = Dimensionless Film Thickness

$h_1$  = Inlet Film Thickness(mm)

$\bar{h}_1$  = Dimensionless Inlet Film Thickness

$h_2$  = Minimum Film Thickness(mm)

$p$  = Film Pressure( $\text{N}/\text{mm}^2$ )

$p_a$  = Ambient Pressure( $\text{N}/\text{mm}^2$ )

$r$  = Radial Coordinate, radius (mm)

$\bar{r}$  = Non-dimensionalised radius,  $r / R_2$

$R_1$  = Inner Radius of Sector(mm)

$R_2$  = Outer Radius of Sector(mm)

$s$  = pitch of bump foil(mm)

$E$  = Elasticity of the top foil material( $\text{N}/\text{mm}^2$ )

$t$  = Thickness of bump foil(mm)

$l$  = Half Bump Length(mm)

$\mu$  = Co-efficient of Viscosity( $\text{Ns}/\text{mm}^2$ )

$\nu$  = Poisson's Ratio of Bump Foil

$\omega$  = Angular Speed of rotation(RPM)

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## **CHAPTER 1**

### **1.1 INTRODUCTION**

One of the major problems of developing turbo-expander system for gas liquefaction plants is the instability of the rotor at high rotational speed. For stability of rotor system at high rotational speed, better bearings are required.

Gas bearings are one of the solutions of maintain stability and prevent contamination of working fluids. Gas bearings are of various types can be used in miniature turbines like Aerostatic gas bearings (externally pressurized gas bearings) and Aerodynamic gas bearings (self-acting).

Gas lubricated externally pressurized bearings consume process gas and they are suitable only up to a medium rotational speed due to whirl speed limitation. Various types of aerodynamic gas bearings can be used are tilting pad journal bearings, spiral groove thrust and journal bearings etc. The major issue with these aerodynamic gas bearings is inability to damp vibrations due to hard supporting surface. The focus of this project is to analyse thrust foil bearings with high ability to damp vibrations at high rotational speed.

A significant volume of component level research has led to recent acceptance of gas foil bearings in several specialized applications like Micro-turbine generators, high speed electric motors, and electrically driven centrifugal blowers etc. Foil bearing supported turbomachinery can benefit from design simplicity and reduced weight, high speed and reduced maintenance. Foil bearings have proven themselves in relatively small lightly loaded applications, like aircraft air cycle machines (ACM's). Recent advances in foil air bearing design, solid lubrication, and bearing and rotor system analytical modelling enable new applications in Oil-Free turbomachinery.

In the present work compressible Reynolds equation is developed and solved based on Finite difference Analysis (FDA) to predict bearing performance parameters, also the result was compared with the rigid bearings.



## 1.2 Applications of Gas Bearings

The first gas journal bearing was demonstrated by Kingsbury (1897). Gas-lubricated bearings are used in many industrial applications in which the hydrodynamic film of gaseous fluid is produced by hydrodynamic action. The gas is generally air. This avoids the need for a liquid lubrication system, simplifies the bearing design and reduces maintenance. Gas bearings are used in gyroscopes where precision and constant torques are required, machine tool spindles, turbo-machinery, dental drills, food and textile machinery and tape and disk drives as part of magnetic storage devices. Gas bearings are also called aerodynamic or self-acting gas bearings.

So far the special case of liquid-lubricated bearings has been considered because the density of liquids can be assumed to be constant. In gas-lubricated bearings, the gas is incompressible and the change in density as a function of pressure and cannot be neglected in the solution of Reynolds' Equation.

Intense development of gas lubrication technology was triggered by the demands of sophisticated navigation systems by the prospects for gas-cooled nuclear reactors, by the proliferation of magnetic peripheral devices in the computer industry and by the everlasting quest for machinery and devices in the aerospace applications. The advantages of gas lubrication are fully established in the following areas:

- i. Machine Tools – Use of gas lubrication in grinding spindles allows attainment of high speeds with minimal heat generation.
- ii. Metrology – Air bearings are used for precise linear and rotational indexing without vibration and oil contamination.
- iii. Dental Drills – High-speed air-bearing dental drills are now a standard equipment in the profession.
- iv. Airborne air-cycle Turbomachines – Foil-type bearings have been successfully introduced for air-cycle turbomachines on passenger aircraft. Increased reliability, leading to reduced maintenance costs, is the benefit derived from air bearings.
- v. Computer peripheral devices – Air lubrication makes possible high-packing-density magnetic memory devices including tapes, discs and drums. Read-write heads now operate at sub micrometre separation from the magnetic film with practically no risk of damage due to wear.

### **1.3 Types of Air Bearings**

Gas bearings are of various types can be used in miniature turbines like:

- a. Aerostatic gas bearings (Externally pressurized gas/air bearings)
- b. Aerodynamic gas bearings (self acting gas/air bearings).

#### **1.3.1 Aerostatic gas bearings:**

1. Aerostatic gas bearings are also known as Externally Pressurized gas bearings and an external pressurized air or process gas is used to maintain pressure between bearing sleeve and the journal.
2. Aerostatic bearings utilize a thin film of high-pressure air to support a load. Since air has a very low viscosity, bearing gaps need to be small, on the order of 1-10  $\mu\text{m}$ .
3. There are five basic types of aerostatic bearing geometries: single pad, opposed pad, journal, rotary thrust, and conical journal/thrust bearings similar to hydrostatic bearings.

#### **1.3.2 Aerodynamic gas bearings:**

1. Aerodynamic are also known as self-acting bearings and an air film is created by the relative motion of two mating surfaces separated by a small distance. From rest, as the speed increases, a velocity induced pressure gradient is formed across the clearance.
2. The increased pressure between the surfaces creates the load carrying effect. The load capacity is dependent on the relative speed at which the surface moves and therefore at zero speed, the bearing supports no load. Zero loads at zero speed effect causes starting and stopping friction and results in some wearing of the bearing surfaces.
3. Despite some of the disadvantages, self-acting bearings have found widespread use in industry. The magnetic read/write heads in disk memory storage devices are in fact aerodynamic bearings that float in close proximity to the disk. This bearing's principal advantage is its ability to act without an external pressure source and solve vibration issues.

## **1.4 ADVANTAGES OF COMPLIANT THRUST BEARINGS OVER OTHER CONVENTIONAL GAS BEARINGS**

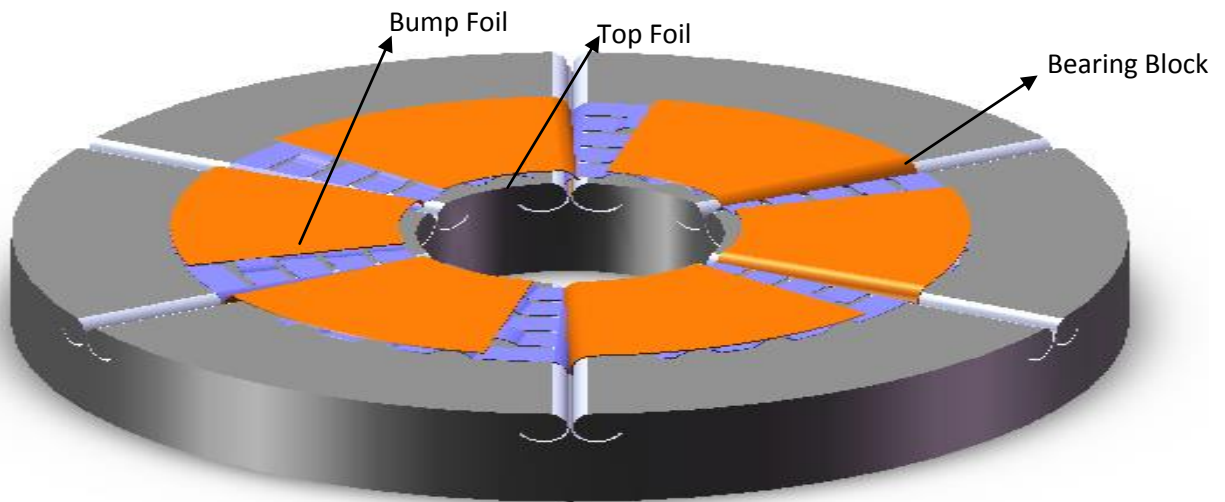
Compliant bearings popularly known as foil bearings have gained very much attention in recent years because of their unique mode of operation and diversity and variety of applications. These types of bearings have various advantages compared to the conventional rigid thrust bearings as they have higher load carrying capacity, less power loss, more stability and greater endurance. These bearings are self-acting and can operate with ambient air or any process gas as the lubricating fluid. The need for complex lubrication systems is eliminated, which results in significant weight reduction and less maintenance. The most common lubricant used is air which is easily available and can operate at high temperatures whereas conventional oil-based lubricants fail since their viscosity drops to a large extent with rise in temperature. These bearings are of importance in the aerospace industry with regard to reduction in weight as well as operating at high speeds and adverse conditions.

## **1.5 COMPLIANT FOIL THRUST BEARINGS**

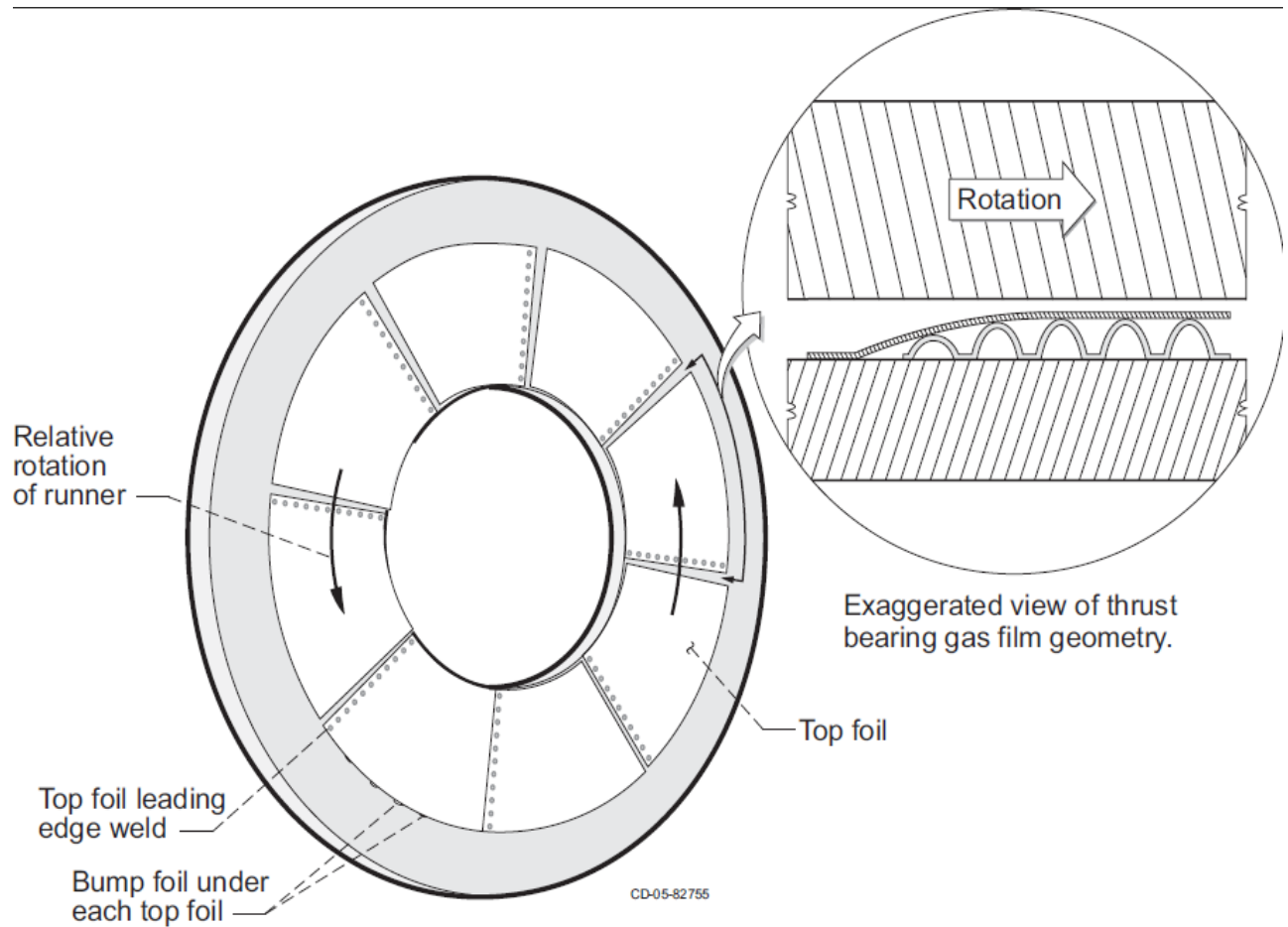
It is consisted of an outer bearing sleeve or outer housing which houses the series of bumps on a thin foil strip and over the bump foil strip a thin smooth top foil sheet is laid (Figure 1.1, 1.2 & 1.3). These foils are welded at one end (starting edge) and are free at the other (trailing edge). The series of bumps in the strip supports the top foil sheet and acts as a spring bed which makes the bearing compliant. The thrust pad and the foils are in contact when the shaft is stationary and remain in contact until a critical lift-off speed(threshold speed) is achieved at which point the thrust rides on a thin gas film developed due to the hydrodynamic pressure developed between thrust pad and the bearing. Due to the action of pressure, the top foil starts to deform. The bump foil is present below the top foil, as the top foil tends to deform, the top foil touches the bump foils and the bump foils act as a spring having some stiffness providing support to the load carried. In compliant thrust bearings, the load direction is parallel to the axis of the rotation of the shaft. During normal operation of the foil bearing which supports the machine, the rotation of the rotor generates a pressurized gas film which pushes the top foil upwards and separates the top foil from the surface of the rotating shaft. The pressure in the air film is proportional to the relative

surface velocity between the rotor and top foil. Thus, faster the rotor rotates, the higher the pressure, and the more is the load the bearing can support. In addition, micro-sliding between top foil and bump foil generates coulomb damping which can increase the dynamic stability of the rotor-bearing system. The dynamic behaviour of a rotating system is significantly influenced by the structural (stiffness and damping) characteristics of the bearings. The exact values of the stiffness and damping coefficients of air-foil bearings are difficult to predict.

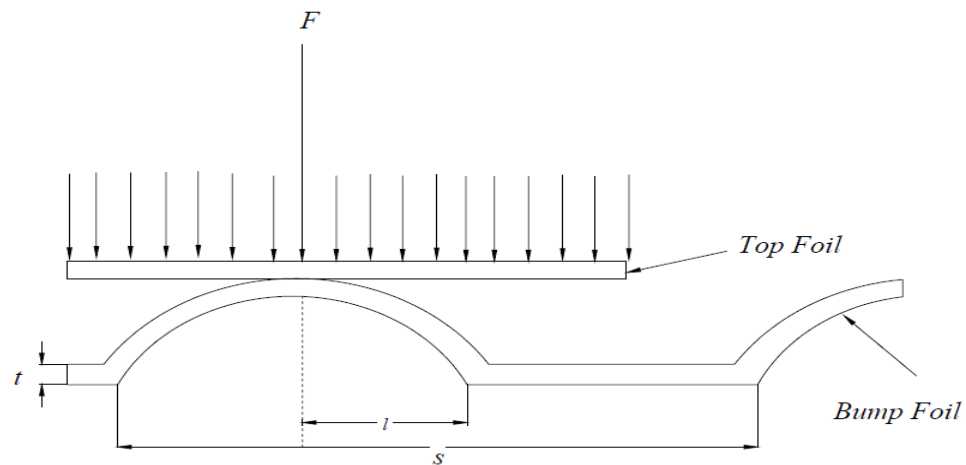
The hydrodynamic pressure developed varies with operating speed and has a significant influence on the deformation of the foils. Hence, the film thickness is a function of hydrodynamic pressure and the elastic properties of the foils. The film thickness is the gap between the top foil and the thrust pad of the shaft. The film thickness is varying with rotation of the shaft and also the pressure exerted on it.



**Fig 1.1 Compliant Foil Thrust Bearing**



**Fig.1.2 Figure depicting construction of Thrust bearings describing foils(top foil and bump foil)**  
 (Source : Heshmat, HWalowit, JA., Pinkus, O,2005)



**Fig 1.3 Arrangement of top and bump foil ( Source : Heshmat, HWalowit, JA., Pinkus, O,2005)**

### **LITERATURE SURVEY**

**Crystal A. Heshmat, David S.Xu and Hooshang Heshmat**[1] investigated analytically utilizing a contemporary approach which combines Finite Difference (FD) and Finite Element (FE) methods. Solution of the governing hydrodynamic equations dealing with compressible fluid is coupled with the structural resiliency of the foil bearing surfaces. FD method is used for hydrodynamic analysis while FE method is used to model structural resiliency. The solution of the hydrodynamic equations using influence coefficients was obtained by using approximations of FD and FE method. Within 2 to 3 iterations, the convergence condition was reached. The overall logic of the program/code written proved to be an efficient technique to deal with the complex structural compliance of various foil bearings.

**Ku and Heshmat** [2] presented an analytical model predicting the deformation of a bump foil. The total elastic bending moment within the bump foils was evaluated by treating the friction forces as conservative forces. The model predicted a higher stiffness if a frictional force is introduced between the bumps and the bearing sleeve and/or the top foil. The bumps near the fixed end have a higher stiffness than the bumps near the free end. In the following paper, Ku and Heshmat examined the bump deflection to verify their analytical model, as presented in [2]. They used an optical tracking device to measure the bump deflection and the recorded hysteresis loops of the bump foil indicated Coulomb damping between the bump foil and the contacting surfaces. In addition, they presented the bump stiffness versus applied loads curves in order to show the effect of bump geometry on stiffness.

**Peng and Carpino** [3] adopted the bending moment equation given in [3] and used an energy method to calculate slip and equivalent viscous damping coefficients under conditions of pure sinusoidal excitation. Peng and Carpino [3] also predicted the stiffness and damping coefficients of foil gas bearings, employing a perturbation and finite element methods. In the analysis of foil bearings, stiffness which was predicted increases with the rotor speed and decreases with the compliance of the bump foil. At low speeds, the overall bearing stiffness depends on the generated hydrodynamic pressure. However, at high speeds, the overall bearing stiffness relies on the structural stiffness of the bump foil, because the

stiffness of the hydrodynamic film is very high. The effect of bearing number on the dynamic force coefficients is presented as well.

**Iordanoff [4]** introduced a simple method to design foil gas thrust bearings. He used a composite profile consisting of a constant slope in the leading edge and a parallel surface to the bearing runner.

**R G Chen, Q Zhou, Y Liu and Y Hou [5]** first proposed a simple type of aerodynamic foil thrust bearing with an elastic hemispherical convex dot support configuration. Then the experimental procedures were carried out on stability and its load capacity characteristics for this foil thrust bearing were conducted on a multi-functional thrust bearing test rig. The preliminary measurement and analysis are presented through the wave and spectrum of axial displacement response in the time and frequency domain.

**Joseph Robert Dickman [6]** manufactured three identical open source foil gas thrust bearings and carried out tests from 0-40,000 RPM against two PS400 coated runners and the results are recorded. He found that bearing torque is proportional to speed and increases linearly with load once the bearing is fully lifted off. Load capacity increases linearly with speed until thermal effects cause top foil distortions and thermal runaway. Beam deflection is compared to the bearing performance.

**Steve Bauman [7], NASA** described in this paper a new test apparatus capable of testing thrust foil air bearings up to 100 mm in diameter at speeds up to 80,000 rpm and temperatures to 650 °C. Bearing Torque, load capacity and bearing temperatures can be measured using this test rig. A number of thrust bearings were tested and intentionally failed with no resultant damage to the test rig. Several test conditions like specific speeds and loads showed about undesirable axial shaft vibrations which have been attributed to the magnetic bearing control system and are under observation. This test rig will be a valuable tool for thrust foil bearing research, parametric studies and technology development.

The work carried out by **Luis San Andres, Tae Ho Kim [8]** performance of Gas Foil Bearings depends largely on the support elastic structure i.e. a smooth foil on top of bump foils. More complex finite element (FE) models couple the elastic deformations of the 2D shell or 1D beam-like top foil to the

bump deflections as well as to the gas film hydrodynamics. It is found that for a 2D FE model predictions overestimate the minimum film thickness at the bearing centreline, while underestimating at the bearing edges. Test data produce the experimental wavy-like film thickness profile. Experimental results of stiffness and damping coefficients versus excitation frequency show that the two FE models result in slightly lower direct stiffness and damping coefficients than those from the simple elastic foundation model.

The work done in this paper by **Robert Bruckner, Brian Dykas, Joseph Prah** [9] describes the methodology for the design and construction of simple foil thrust bearings for performance testing and low marginal costs is presented. The design of fixtures and tooling required to fabricate foil thrust bearings is presented, using conventional machining processes where possible. A prototype bearing is constructed using all the steps required for fabrication. A load-deflection curve for the bearing is presented to illustrate structural stiffness characteristics. The results are that performance of these bearings seemed to be useful when compared with data from the open literature or conventional rigid bearings.

The authors **Quan Zhou, Yu Hou, Chunzheng Chen** [10] carried static and stability experiments on a high speed turbine test rig. This paper describes the results of a gas thrust bearing with viscoelastic support which is meant for high speed turbo-machinery. The gas bearing, which belongs to compliant foil bearings, consists of a top thin metal foil and a bottom thin rubber foil. The static results show that the structural stiffness of test bearing generally increases with the increase in axial load and the decrease in thickness of bottom foil. In the rotation tests, rotor runs stably with small vibration amplitude, which is dominant in waterfall plot during whole speed up procedure. The results show that the bearing has good stability characteristics for high speed gas turbines.

The experiment carried out in this paper by **Brian David Dykas**[11] is about studying the operating characteristics of foil gas thrust bearings experimentally and analytically to know more about the physical mechanisms that limit the bearing performance. Variety of configurations were used to measure bearing power loss and load capacity which highlight several important factors that influence



the performance. It is found that according to the conventional hydrodynamic theory, surface condition of the foil and surface condition of the runner have a large influence on bearing performance. Thermal effects are found out to be more at higher loads where gas film heat generation and resulting thermoelastic distortion are larger, but smooth surfaces with lubrication are needed to achieve these loads. This paper in general summarises the effects of these non-ideal surface conditions on the load capacity of foil thrust bearings.

The work done by scientists **Vikas Arora, P.J.M. van der Hoogt and R.G.K.M. Aarts[12]** describes about the experimental procedures to identify the stiffness and damping characteristics of Axial Air Foil Bearings. The innermost (top foil) traps a gas pressure film that supports a load while the layers below provide an elastic foundation. Identification of structural characteristics is important for successful design practice. Experiments are carried at a maximum speed of 60,000 rpm. Sub-structuring approach is used for identification of the structural i.e. stiffness and damping characteristics of the Air Foil Bearings.

The paper by **Robert J. Bruckner, NASA [13]** illustrates the experimental results of the performance of simple gas foil thrust bearing in air. This experiment is carried out to provide machine designers the basic performance parameters and to explain the underlying physics of foil thrust bearings. The tests were conducted on simple bump foil supported thrust bearings. Test conditions consist of air at ambient pressure and temperatures up to 500<sup>0</sup>C and speeds to 55,000 rpm. A complete set of axial load, frictional torque and rpm is obtained for two different compliant sub-structures and inter-pad gaps. Data obtained from commercially available foil thrust bearings both with and without active cooling is presented for comparison. Speed-load characteristic of foil thrust bearing is found out from the test results. For the journal bearing, the load capacity increases linearly with rpm but for the thrust bearing operates in the hydrodynamic high speed limit.

In this paper, a generalized hydrodynamic analysis is carried out by **Robert Jack Bruckner [14]** to analyse individual effects included in the development of the governing equations. The governing equations are the conservation equations of mass, momentum and energy. The Reynolds' equation is

developed from these equations. The energy equation is simplified by applying the thin film layer assumptions such that fluid properties do not vary through the film. The structural deformation of the bearing is modelled with a single partial differential equation. A linear superposition of hydrodynamic load and compliant foundation reaction is reached. The stiffness of the compliant foundation can be seen as a set of springs that support the top foil. This system of governing equations is solved by using numerical methods by writing a computer program in Mathematics computing environment. This work finds a substantial difference between bearing performance based on traditional lubricant models and that based on the energy equation model.

The article by **Yong-Bok Lee, Tae Young Kim, Chang Ho Kim and Tae Ho Kim [15]** explains about a model of thrust bump foil bearings which predicts the deflection with variable axial load with an assumption that there is no tilting effect of the thrust collar. To predict the air clearance, deflection of the elastic foundation was used in the air film height equation. Combined Dirichlet and Neumann-type boundary conditions were used for static load performance predictions. To verify the theoretical data and conditions, experiments were carried out with three different thrust foil bearings. The rpm was varied for all the three bearings. It was found out that the model using nonlinear stiffness was in better agreement with the experimental results than the model using linear stiffness.

### **MATHEMATICAL MODEL**

#### **3.1 DERIVATION OF STANDARD REYNOLDS' EQUATION FOR THRUST BEARINGS**

The most accurate way of predicting the performance parameters for any type of thrust bearings is to solve the lubricant flow equations obtained from Navier-Stokes relationships. The solutions become very complex and the costs for computation become very high when variations in viscosity and flexibility of outer bearing sleeves are considered.

Hence instead of full solution, some approximate methods are used for solving the two dimensional Reynolds' Equation. Reynolds' Equation gives tremendous insight into fluid behaviour in bearing lubricant films, and forms the basis for birth of science of hydrodynamic lubrication. Solution of the traditional Reynolds' Equation helps us to determine the pressure distribution in a bearing having a random film shape. After the pressure profile is evaluated the important bearing parameters such as load-carrying capacity, friction force, flow rates, etc. can be easily determined.

The standard Reynolds' Equation is considered by neglecting the time variant. The solution to this equation is obtained by considering a bearing with finite length and other dimensions. This approximation represents the most accurate method to predict the performance parameters for any type of thrust bearings.

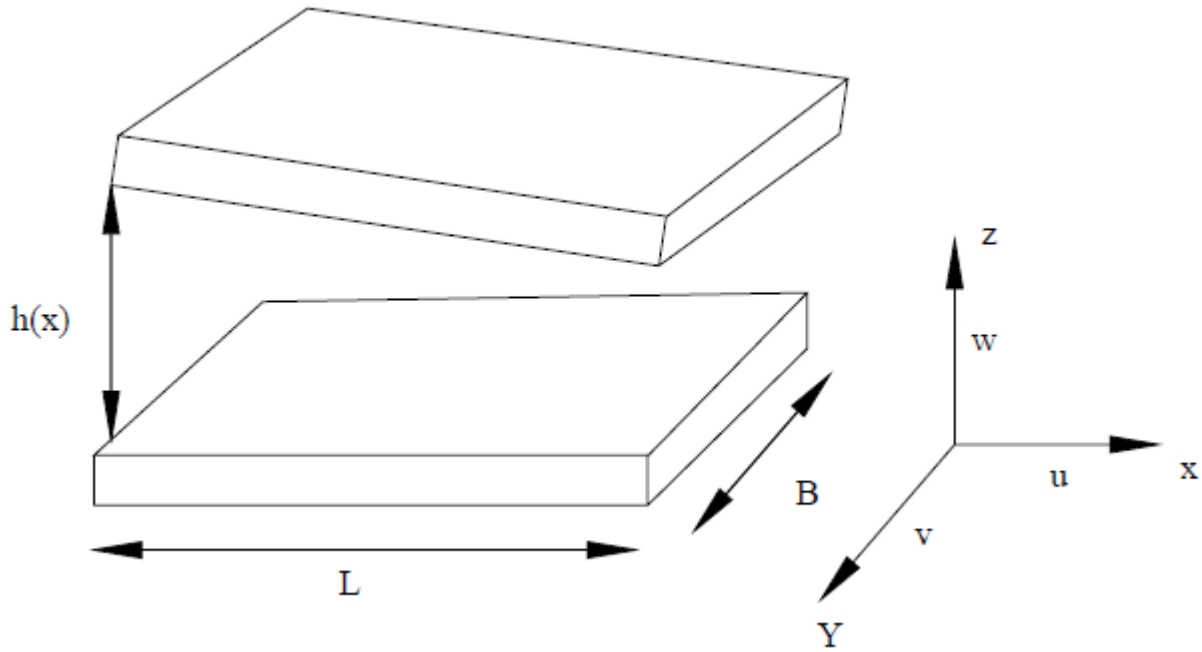
#### **3.2 Assumptions taken in Derivation of Reynolds Equation**

The following assumptions are made in deriving the Reynolds Equation

1. Fluid is assumed Newtonian, with direct proportionality between shear stress and Shearing velocity  $\tau \propto \frac{dv}{dy}$
2. Inertia and body forces are assumed to be negligible compared to the viscous terms.
3. Variation of pressure across the film is assumed to be very small.
4. Flow is laminar.
5. Curvature effects are negligible.

### 3.3 Compressible Reynolds' equation in two dimensions

#### 3.3.1 Use of Conservation of Momentum



**Fig. 3.1 :Schematic Of Slider Bearing(Source:Thesis By Hanumappa Reddy)**

Referring to figure 3.1

Let  $h=h(x,y)$  be the film gap between runner and to top foil.

X- Direction Momentum Equation

$$\frac{\partial p}{\partial x} = \frac{\partial}{\partial z} \left( \mu \frac{\partial u}{\partial z} \right) \quad \text{-----} \quad (3.1)$$

Y-Direction Momentum Equation

$$\frac{\partial p}{\partial y} = \frac{\partial}{\partial z} \left( \mu \frac{\partial v}{\partial z} \right) \quad \text{-----} \quad (3.2)$$

Z Direction Momentum Equation

$$\frac{\partial p}{\partial z} = 0 \quad (\text{Because we have assumed that variation of pressure across the film is very small})$$

Integrating 3.1., we get,

$$\left(\frac{\partial p}{\partial x}\right)z = \mu\left(\frac{\partial u}{\partial z}\right) + C_1$$

$$\Rightarrow \left(\frac{\partial p}{\partial x}\right)\frac{z^2}{2} = \mu u + C_1 z + C_2$$

Using Boundary condition at  $z=0$ ,  $u=u_a$

$$0 = \mu u_a + C_2$$

$$\Rightarrow C_2 = -\mu u_a$$

Similarly at  $z=h$ ,  $u=u_b$

$$\left(\frac{\partial p}{\partial x}\right)\frac{h^2}{2} = \mu u_b + C_1 h - \mu u_a$$

$$\Rightarrow C_1 = \left(\frac{\partial p}{\partial x}\right)\frac{h}{2} + \frac{\mu}{h}(u_a - u_b)$$

Now putting values of  $C_1$  and  $C_2$ , we get

$$\mu u = \left(\frac{\partial p}{\partial x}\right)\frac{z^2}{2} - C_1 z - C_2$$

$$\Rightarrow \mu u = \left(\frac{\partial p}{\partial x}\right)\frac{z^2}{2} - \left(\frac{\partial p}{\partial x}\right)\frac{h}{2}z - \frac{\mu z}{h}(u_a - u_b) + \mu u_a$$

$$\Rightarrow u = \frac{1}{2\mu}\left(\frac{\partial p}{\partial x}\right)\frac{z^2}{2} - \left(\frac{\partial p}{\partial x}\right)\frac{hz}{2\mu} - \frac{z}{h}(u_a - u_b) + u_a$$

$$\Rightarrow u = \frac{1}{2\mu}\left(\frac{\partial p}{\partial x}\right)(z^2 - zh) + u_a\left(1 - \frac{z}{h}\right) + \frac{z}{h}u_b$$

Now integrating 3.2. We get

$$\left(\frac{\partial p}{\partial y}\right)z = \mu\left(\frac{\partial v}{\partial z}\right) + C_1$$

$$\Rightarrow \left(\frac{\partial p}{\partial y}\right)\frac{z^2}{2} = \mu v + C_1 z + C_2$$

At  $z=0$ ,  $v=0$

Which gives  $C_2 = 0$

At  $z=h$ ,  $v=0$

$$\left(\frac{\partial p}{\partial y}\right)\frac{h^2}{2} = C_1 h$$

This gives  $C_1 = \left(\frac{\partial p}{\partial y}\right)\frac{h}{2}$

Now putting values of  $C_1$  and  $C_2$ , we get

$$\left(\frac{\partial p}{\partial y}\right) \frac{z^2}{2} = \mu v + \left(\frac{\partial p}{\partial y}\right) \frac{hz}{2}$$

$$v = \frac{1}{2\mu} \left(\frac{\partial p}{\partial y}\right) (z^2 - zh) \text{----- (3.3)}$$

With X being the direction of sliding and Y being the direction of leakage of fluid

### 3.3.2 Use of Conservation of Mass

Applying conservation of mass,

$$\int_0^h \frac{\partial}{\partial t} (\rho) dz + \int_0^h \frac{\partial}{\partial x} (\rho u) dz + \int_0^h \frac{\partial}{\partial y} (\rho v) dz + \int_0^h \frac{\partial}{\partial z} (\rho w) dz = 0 \text{----- (3.4)}$$

The final equation comes out to be:

$$h \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} \left[ -\frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial x}\right) + \frac{\rho h}{2} (u_a + u_b) \right] + \frac{\partial}{\partial y} \left[ -\frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial y}\right) \right] + \rho (w_b - w_a) = 0$$

$$\Rightarrow \frac{\partial}{\partial x} \left( \frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial x}\right) \right) + \frac{\partial}{\partial y} \left( \frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial y}\right) \right) + h \frac{\partial \rho}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} (\rho h (u_a + u_b)) + \rho (w_b - w_a)$$

For a journal bearing, bearing sleeve is fixed and the journal is rotating. Hence  $u_a=0$  and  $w_a=0$ .

Also  $u_b=u$  is not 0 and  $w_b=w$  is not 0. Also neglecting the time variant and considering the above equation in Two dimensions:

$$\frac{\partial}{\partial x} \left( \frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial x}\right) \right) + \frac{\partial}{\partial y} \left( \frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial y}\right) \right) = \frac{1}{2} \frac{\partial}{\partial x} (\rho h (0 + u))$$

$$\frac{\partial}{\partial x} \left( \frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial x}\right) \right) + \frac{\partial}{\partial y} \left( \frac{\rho h^3}{12\mu} \left(\frac{\partial p}{\partial y}\right) \right) = \frac{1}{2} u \frac{\partial}{\partial x} (\rho h) \text{----- (3.5)}$$

### 3.3.3 Reynolds Equation in Polar Coordinates

For thrust bearing expression for Reynolds Equation in Polar Coordinates is required,

$$x = R\theta \Rightarrow \partial x = R\partial\theta$$

$$u = R\omega$$

$$\Rightarrow \frac{\partial}{R\partial\theta} \left( \frac{\rho h^3}{12\mu R} \left( \frac{\partial p}{\partial\theta} \right) \right) + \frac{\partial}{\partial y} \left( \frac{\rho h^3}{12\mu} \left( \frac{\partial p}{\partial y} \right) \right) = \frac{1}{2} R\omega \frac{\partial}{R\partial\theta} (\rho h) \text{-----} (3.6)$$

Considering ideal gas with  $R_g$  as gas constant i.e.

$$\rho = \frac{p}{R_g T}$$

$$\Rightarrow \frac{\partial}{R^2\partial\theta} \left( \frac{ph^3}{12\mu R_g T} \left( \frac{\partial p}{\partial\theta} \right) \right) + \frac{\partial}{\partial y} \left( \frac{ph^3}{12\mu R_g T} \left( \frac{\partial p}{\partial y} \right) \right) = \frac{1}{2} \frac{\omega}{R_g T} \frac{\partial}{\partial\theta} (ph)$$

Multiplying by  $(R^2 R_g)$ , we get

$$\frac{\partial}{\partial\theta} \left( \frac{ph^3}{12\mu T} \left( \frac{\partial p}{\partial\theta} \right) \right) + R^2 \frac{\partial}{\partial y} \left( \frac{ph^3}{12\mu T} \left( \frac{\partial p}{\partial y} \right) \right) = \frac{\omega R^2}{2} \frac{\partial}{\partial\theta} \left( \frac{ph}{T} \right)$$

Normalizing the above equation with:

$$\bar{y} = \frac{y}{\frac{L}{2}}; \bar{p} = \frac{p}{p_a}; \bar{h} = \frac{h}{C}; \bar{\mu} = \frac{\mu}{\mu_0}; \bar{T} = \frac{T}{T_0}$$

$$\frac{\partial}{\partial\theta} \left( \frac{\bar{p} p_a \bar{h}^3 C^3}{12\bar{\mu} \mu_0 \bar{T} T_0} \left( \frac{\partial}{\partial\theta} (\bar{p} p_a) \right) \right) + R^2 \frac{\partial}{\frac{L}{2} \partial \bar{y}} \left( \frac{\bar{p} p_a \bar{h}^3 C^3}{12\bar{\mu} \mu_0 \bar{T} T_0} \left( \frac{\partial}{\frac{L}{2} \partial \bar{y}} (\bar{p} p_a) \right) \right) = \frac{\omega R^2}{2} \frac{\partial}{\partial\theta} \left( \frac{\bar{p} p_a \bar{h} C}{\bar{T} T_0} \right)$$

$$\left[ \frac{p_a^2 C^3}{12\mu_0 T_0} \right] \frac{\partial}{\partial\theta} \left( \frac{\bar{p} \bar{h}^3}{\bar{\mu} \bar{T}} \left( \frac{\partial}{\partial\theta} (\bar{p}) \right) \right) + \frac{R^2}{\left( \frac{L}{2} \right)^2} \left[ \frac{p_a^2 C^3}{12\mu_0 T_0} \right] \frac{\partial}{\partial \bar{y}} \left( \frac{\bar{p} \bar{h}^3}{\bar{\mu} \bar{T}} \frac{\partial}{\partial \bar{y}} (\bar{p}) \right) = \frac{\omega R^2}{2} \left( \frac{p_a C}{T_0} \right) \frac{\partial}{\partial\theta} \left( \frac{\bar{p} \bar{h}}{\bar{T}} \right)$$

ASSUMPTION: Using Isothermal and Isoviscous case

$$\Rightarrow \bar{\mu} = 1 \text{ and } \bar{T} = 1$$

$$\Rightarrow \frac{\partial}{\partial\theta} (\bar{p} \bar{h}^3 \frac{\partial \bar{p}}{\partial\theta}) + \left( \frac{2R}{L} \right)^2 \frac{\partial}{\partial \bar{y}} (\bar{p} \bar{h}^3 \frac{\partial \bar{p}}{\partial \bar{y}}) = \frac{\omega R^2}{2} * \left( \frac{p_a C}{T_0} \right) * \frac{12\mu_0 T_0}{p_a^2 C^3} \frac{\partial}{\partial\theta} (\bar{p} \bar{h})$$

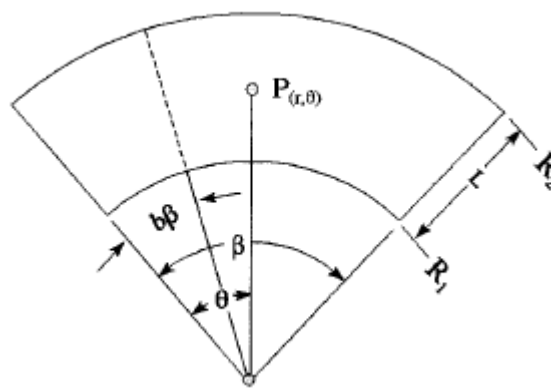
Taking Compressibility or Bearing Number as:

$$\Lambda = \frac{6\omega\mu_0}{p_a} \left(\frac{R}{C}\right)^2$$

We get the standard compressible Reynolds Equation in 2D as:

$$\frac{\partial}{\partial \theta} (\bar{p} \bar{h}^3 \frac{\partial \bar{p}}{\partial \theta}) + \left(\frac{D}{L}\right)^2 \frac{\partial}{\partial \bar{y}} (\bar{p} \bar{h}^3 \frac{\partial \bar{p}}{\partial \bar{y}}) = \Lambda \frac{\partial}{\partial \theta} (\bar{p} \bar{h}) \text{ -----(3.7)}$$

### 3.4 Single Pad of Thrust Bearings



**Fig. 3.2 :Schematic of Single Pad of Thrust Bearing(Source:Crystal A. heshmat, David Xu(2000),Journal of tribology,Vol.122)**

Using Dimensionless parameters:

$$\bar{h} = \frac{h}{h_2}, \bar{r} = \frac{r}{R_2}, \bar{g} = \frac{g}{h_2}, \bar{h}_1 = \frac{h_1}{h_2}$$

And using the Bearing Compression factor (Bearing Number):

$$\Lambda = \frac{6\mu\omega}{p_a} * \left(\frac{R_2}{h_2}\right)^2$$

The Dimensionless Reynolds Equation becomes:

$$\frac{1}{\bar{r}} \frac{\partial}{\partial \bar{r}} \left( \bar{r} \bar{h}^3 \cdot \bar{p} \cdot \frac{\partial \bar{p}}{\partial \bar{r}} \right) + \frac{1}{\bar{r}^2} \frac{\partial}{\partial \theta} \left( \bar{h}^3 \cdot \bar{p} \cdot \frac{\partial \bar{p}}{\partial \theta} \right) = \Lambda \frac{\partial (\bar{p} \cdot \bar{h})}{\partial \theta}$$



The film thickness is a function of wedge shape geometry and pressure at each point and expressed as :

$$\bar{h} = 1 + \bar{g}(\bar{r}, \theta) + \alpha(\bar{p} - 1)$$

Where

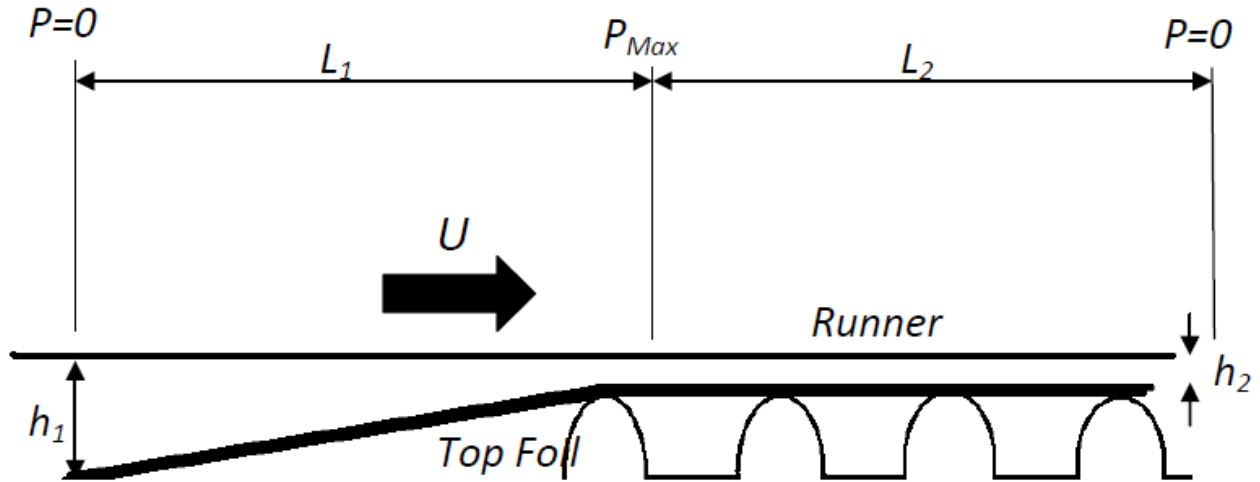
$$\bar{g} = (\bar{h}_1 - 1) \left( 1 - \frac{\theta}{b\beta} \right) \text{ when } 0 \leq \theta \leq b\beta$$

$$= 0 \text{ when } b\beta \leq \theta \leq \beta$$

where b is the extent of the top foil

$$\alpha = \frac{2p_a s}{cE} \left( \frac{l}{t} \right)^3 (1 - \nu^2)$$

$\alpha$  is the compliance coefficient.



**Fig 3.3 Variation of film thickness and Boundary Conditions (Source: Heshmat, Xu (2000), Vol.122 )**

$$p = p_a \text{ at } r = R_1 \text{ and } R_2$$

$$p = p_a \text{ at } \theta = 0 \text{ and } \beta$$

These are the boundary conditions which is to be used to find out the solution of Reynolds' Equation

**NUMERICAL METHODS**

Some approximate numerical methods must be adopted to solve the Reynolds' Equation. In order to find the pressure profile, we are using FINITE DIFFERENCE METHOD. This chapter consists of a detailed description of the Finite Difference Method which is used for discretisation of the Reynolds' Equation. The Reynolds' Equations was written in Finite Difference form and solved by means of an iterative procedure. The domain i.e. radius(R) and theta is divided into segments of small size and a mesh was generated. The pressure at each node is obtained by iterative procedure.

The Reynolds' Equation is as follows

$$\frac{1}{r} \frac{\partial}{\partial r} \left( \frac{\bar{r}^3}{r \cdot \bar{h}} \cdot \bar{p} \cdot \frac{\partial \bar{p}}{\partial r} \right) + \frac{1}{r^2} \frac{\partial}{\partial \theta} \left( \frac{\bar{h}^3}{\bar{h}} \cdot \bar{p} \cdot \frac{\partial \bar{p}}{\partial \theta} \right) = \Lambda \frac{\partial (\bar{p} \cdot \bar{h})}{\partial \theta} \quad \text{----- (4.1)}$$

**4.1 FINITE DIFFERENCE FORMS**

From central finite difference form we get

$$\begin{aligned} \frac{\partial \bar{p}}{\partial r} &= \frac{\bar{p}_{i+1,j} - \bar{p}_{i-1,j}}{2\Delta r} \\ \frac{\partial^2 \bar{p}}{\partial r^2} &= \frac{\bar{p}_{i+1,j} - 2\bar{p}_{i,j} + \bar{p}_{i-1,j}}{(\Delta r)^2} \\ \frac{\partial \bar{p}}{\partial \theta} &= \frac{\bar{p}_{i,j+1} - \bar{p}_{i,j-1}}{2\Delta \theta} \\ \frac{\partial^2 \bar{p}}{\partial \theta^2} &= \frac{\bar{p}_{i,j+1} - 2\bar{p}_{i,j} + \bar{p}_{i,j-1}}{(\Delta \theta)^2} \end{aligned}$$

Let Equation 4.1 is written in terms of  $A + B = C$

Then,

$$\begin{aligned} A &= \frac{\bar{p}\bar{h}^3}{\bar{r}} \left( \frac{\partial \bar{p}}{\partial \bar{r}} \right) + \bar{h}^3 \left( \frac{\partial \bar{p}}{\partial \bar{r}} \right)^2 + \bar{p} \left[ 3\bar{h}^2 \left( \frac{\partial \bar{h}}{\partial \bar{r}} \right) + \bar{h}^3 \frac{\partial^2 \bar{p}}{\partial \bar{r}^2} \right] \\ &= \frac{\bar{p}\bar{h}^3}{\bar{r}} \left( \frac{\partial \bar{p}}{\partial \bar{r}} \right) + \bar{h}^3 \left( \frac{\partial \bar{p}}{\partial \bar{r}} \right)^2 + \bar{p}\bar{h}^3 \left( \frac{\partial^2 \bar{p}}{\partial \bar{r}^2} \right) \end{aligned}$$

So writing this differential equation in finite difference form we get

$$A = \frac{\bar{p}_{i,j} \bar{h}_{i,j}^3}{\bar{r}_{i,j}} \left[ \frac{\bar{p}_{i+1,j} - \bar{p}_{i-1,j}}{2\Delta \bar{r}} \right] + \bar{h}_{i,j}^3 \left[ \frac{\bar{p}_{i+1,j} - \bar{p}_{i-1,j}}{2\Delta \bar{r}} \right]^2 + \bar{p}_{i,j} \bar{h}_{i,j}^3 \left[ \frac{\bar{p}_{i+1,j} - 2\bar{p}_{i,j} + \bar{p}_{i-1,j}}{(\Delta \bar{r})^2} \right]$$

Similarly

$$\begin{aligned} B &= \frac{1}{\bar{r}^2} \bar{h}^3 \left( \frac{\partial \bar{p}}{\partial \theta} \right) + \frac{\bar{p}}{\bar{r}^2} \frac{\partial}{\partial \theta} \left( \bar{h}^3 \left( \frac{\partial \bar{p}}{\partial \theta} \right) \right) \\ &= \frac{1}{\bar{r}^2} \bar{h}^3 \left( \frac{\partial \bar{p}}{\partial \theta} \right) + \frac{3\bar{p}\bar{h}^2}{\bar{r}^2} \left( \frac{\partial \bar{h}}{\partial \theta} \right) + \bar{h}^3 \left( \frac{\partial^2 \bar{p}}{\partial \theta^2} \right) \end{aligned}$$

Writing these differential terms in finite difference form we get

$$B = \frac{1}{\bar{r}_{i,j}^2} \bar{h}_{i,j}^3 \left[ \frac{\bar{p}_{i,j+1} - \bar{p}_{i,j-1}}{2\Delta \theta} \right] + \frac{3\bar{p}_{i,j} \bar{h}_{i,j}^2}{\bar{r}_{i,j}^2} \left[ \frac{\bar{h}_{i,j+1} - \bar{h}_{i,j-1}}{2\Delta \theta} \right] + \bar{h}_{i,j}^3 \left[ \frac{\bar{p}_{i,j+1} - 2\bar{p}_{i,j} + \bar{p}_{i,j-1}}{(\Delta \theta)^2} \right]$$

$$\begin{aligned} C &= \Lambda \frac{\partial (\bar{p}\bar{h})}{\partial \theta} \\ &= \Lambda \left[ \left( \frac{\partial \bar{p}}{\partial \theta} \right) \bar{h} + \bar{p} \left( \frac{\partial \bar{h}}{\partial \theta} \right) \right] \end{aligned}$$

Writing these differential terms in finite difference form we get

$$C = \Lambda \bar{h}_{i,j} \left[ \frac{\bar{p}_{i,j+1} - \bar{p}_{i,j-1}}{2\Delta \theta} \right] + \Lambda \bar{p}_{i,j} \left[ \frac{\bar{h}_{i,j+1} - \bar{h}_{i,j-1}}{2\Delta \theta} \right]$$

Substituting these values in the Reynolds' Equation and arranging the equation to form a quadratic equation in  $\bar{p}_{i,j}$  of the form we get:

$$D(\bar{p}_{i,j})^2 + E(\bar{p}_{i,j}) + F = 0$$

The solution of the quadratic equation can be given as

$$\bar{p}_{i,j} = \frac{-E \pm \sqrt{E^2 - 4DF}}{2D}$$

Here the nodes along the theta direction are denoted as “i” and the nodes along radial direction are denoted as “j”.

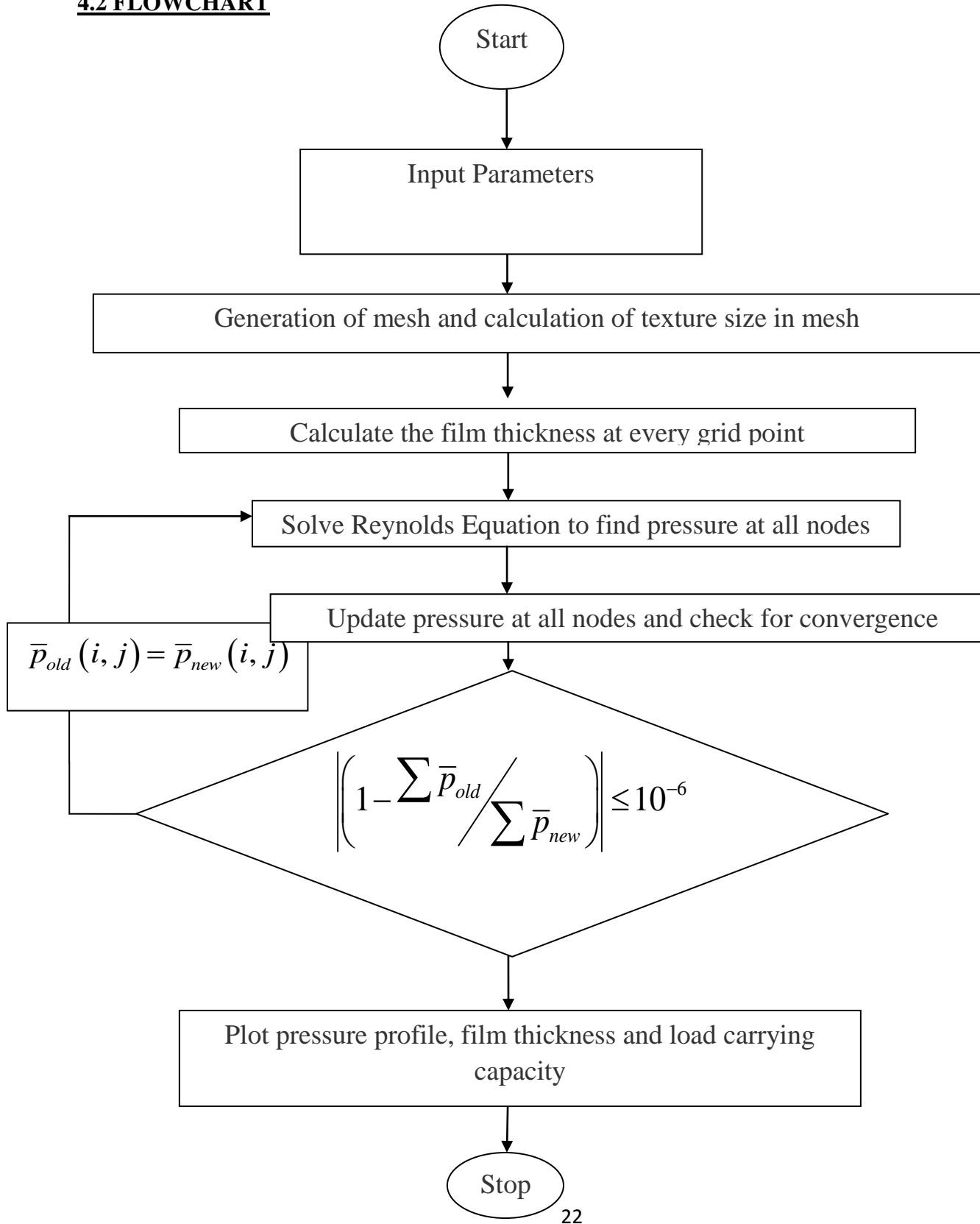
The whole domain can be made as a mesh of ‘i’ and ‘j’ i.e. theta and radius respectively.

To know the pressure at each node, the value of pressure at previous and next of the node both in i<sup>th</sup> and j<sup>th</sup> direction must be known. So, the process is an iterative process. MATLAB (MATRIX LABORATORY) software is used in which the program is written. The program contains loops for ‘i’ and ‘j’ and within the loops the pressure is calculated.

Similarly, the film thickness is also calculated. As film thickness is a function of radius and theta and also pressure at the node, the film thickness is also calculated within the loops. To know the value of film thickness at each node, the value previous and next node must be known. One thing is different here that is the film thickness varies only in theta direction and remains the same along radial direction. We use approximation method to solve this iterative process. At the beginning of the program we assign a certain value of pressure at each and every node. This value assigned is just an arbitrary value. Final values of pressure are calculated at the end of the program when the values of pressure obey a certain convergence condition. Here the convergence condition is values must be same up to 3 decimal places

Initially at the beginning of the program, we calculate the values of compliance coefficient and bearing number. We also give other dimensions of the thrust bearing at the beginning.

## 4.2 FLOWCHART



**RESULTS AND DISCUSSIONS**

The analysis was done for two different types of bump materials keeping the bump geometry same.

The bump foil material chosen are Inconel X 750 and Aluminium bronze. Inconel X-750 is a Nickel-Chromium alloy made precipitation harden able by additions of Aluminium and Titanium, having creep-rupture strength at high temperatures to about 700°C and they are used in high temperature application like gas turbines, rocket engines etc. Whereas Aluminium bronze is applied where high temp is not an issue and it is within 40°C of operating temperature.

The details of input parameters of foil geometry and properties are given in table 5.1.

<b>INPUT PARAMETERS</b>	<b>VALUES</b>
No of grid points in theta direction	50
No of grid points along radial direction	50
Inner radius	10 mm
Outer radius	24 mm
Radial Clearance(=h <sub>2</sub> )	0.03mm
Inlet Film Thickness(=h <sub>1</sub> )	0.2mm
Length of Bearing	23.5mm
Bump pitch	3.17mm
Ambient Pressure	0.1 N/mm <sup>2</sup>
Half bump length	1.125mm
Thickness of Bump Foil	0.1 mm
Poisson's Ratio	0.29
Coefficient of Viscosity	17.8 x 10 <sup>-12</sup> Ns/mm <sup>2</sup>
Angular Speed of the Shaft	1,00,000 RPM
Modulus of Elasticity	160,000 N/mm <sup>2</sup>
No of iterations	300
No of Pads	6

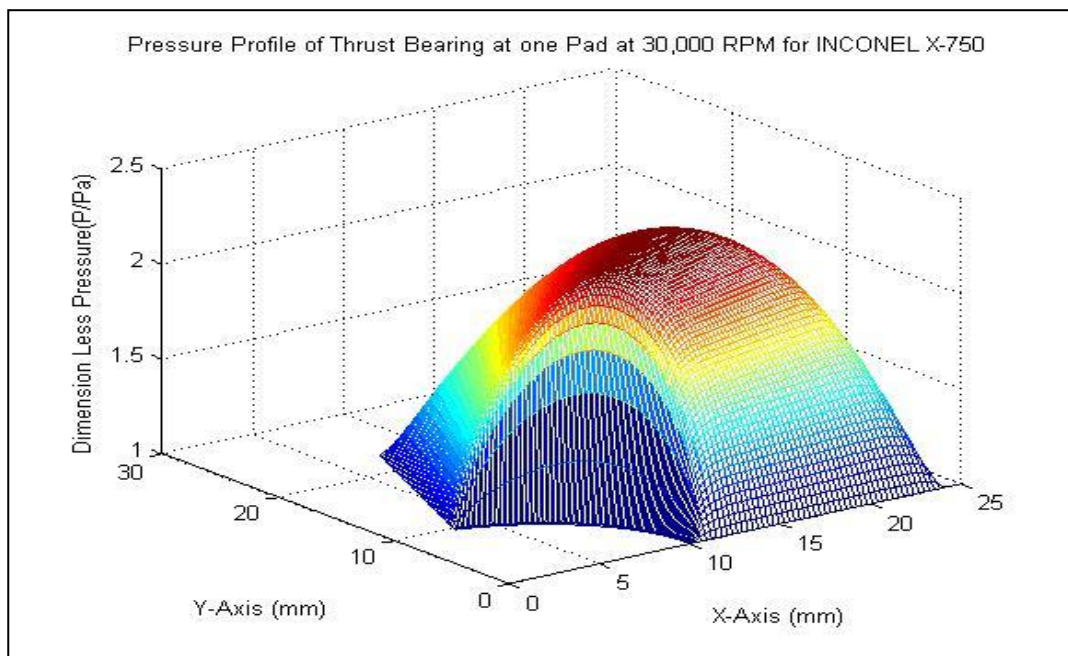
**Table 5.1: Geometry and properties of Bearings parts.**

### 5.1 Analysis with Inconel X-750 foils.

The Modulus of Elasticity used for Inconel X 750 is 212.4 GPa and poisson's ratio is 0.29. Solving the modified Reynolds equation for compliant foil thrust bearings, with the given properties and dimensions from table 1, following analysis are done.

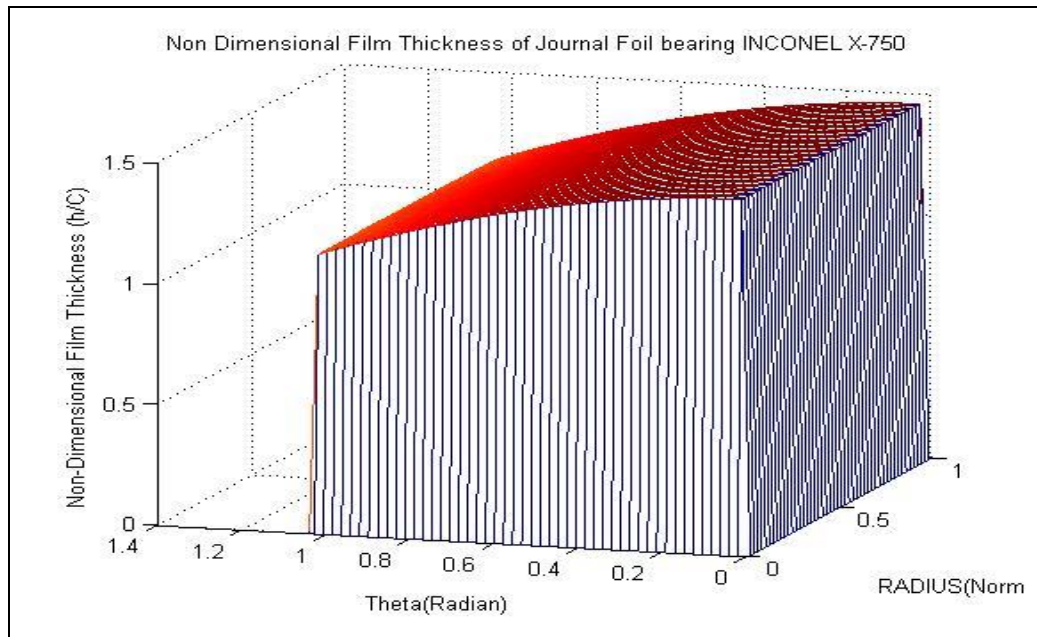
- Pressure profile over single pad.
- Film thickness variation over a single pad
- Comparison of film thickness with rigid bearings.
- Variations of load carrying capacity with variation of speed of runner.
- Variations of load carrying capacity with variations of foil thickness.

Figure 5.1 Shows variation of pressure profile over a single pad of  $60^\circ$ . The maximum dimensionless pressure is about 2.45 and corresponding load carrying capacity at 30000 RPM of runner is 107 N, which is sufficient to carry the axial load of small turbines and turbochargers.



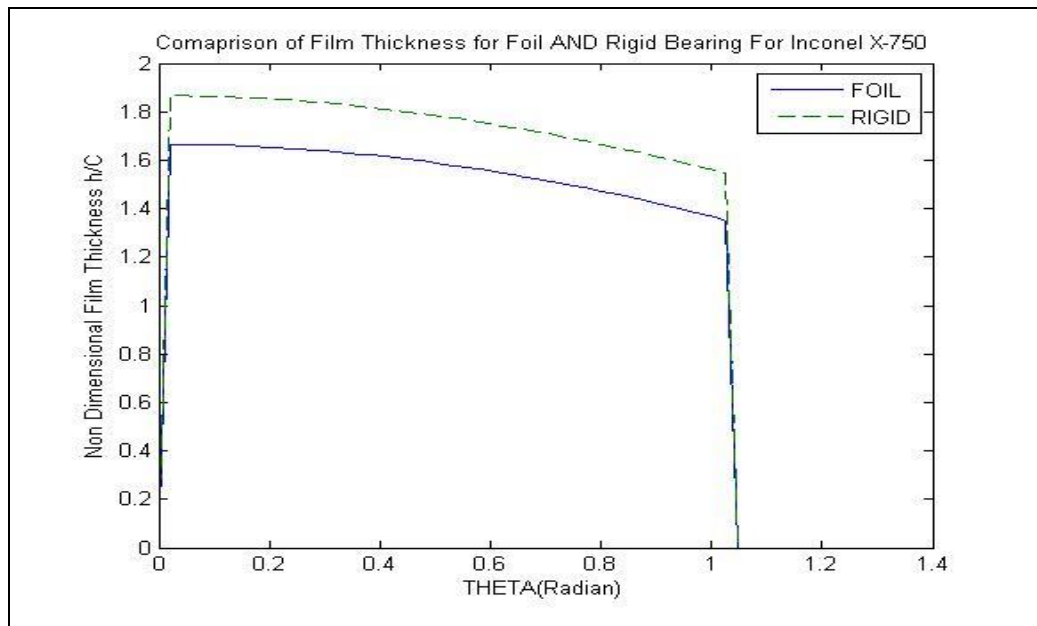
**Fig 5.1: Dimensionless pressure profile over one pad at 30,000 RPM**

Figure 5.2 shows the dimensionless film thickness over one pad. The film thickness maintained is about 200 microns, and this is much above the surface roughness of the bearing surface, which shows a positive sign for aerodynamic bearings.



**Fig 5.2: Dimensionless Film thickness over one pad at 30,000 RPM**

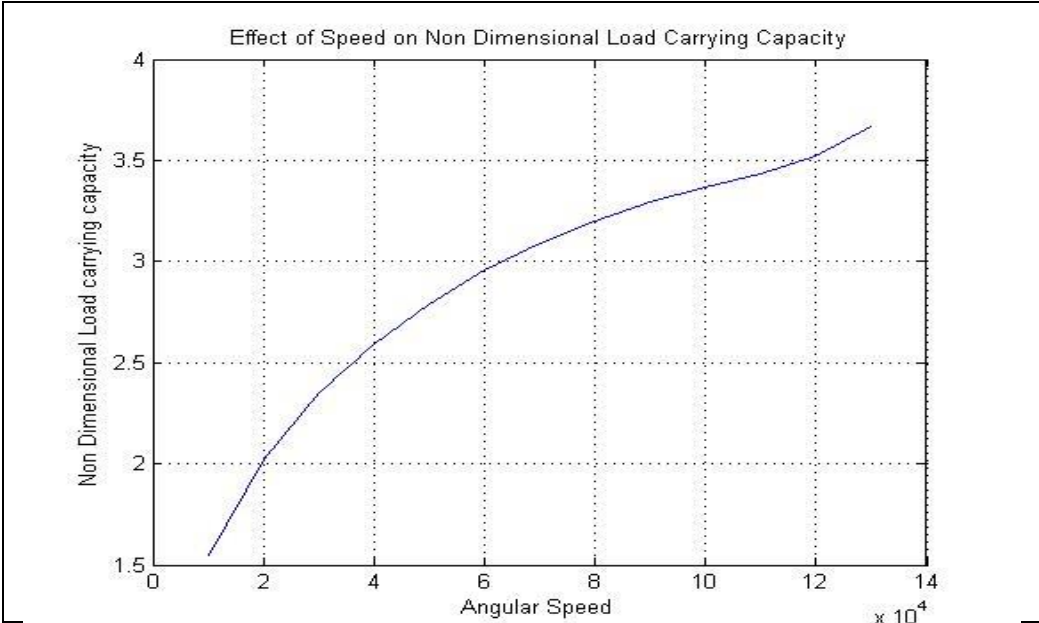
Figure 5.3 shows comparison of film thickness between compliant foil bearings and rigid bearings. The performance of foil bearings seems to be better than rigid bearings.



**Fig 5.3: Comparison of Film thickness between foil and rigid bearings over one pad at 30,000 RPM**

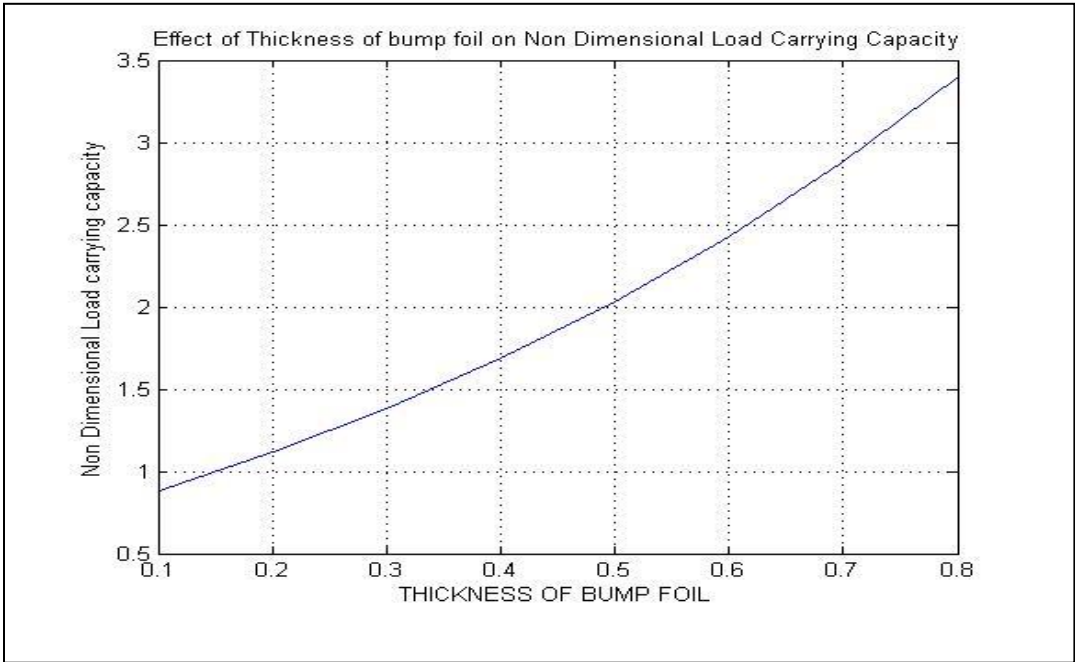


Figure 5.4 shows the variation of load carrying capacity with increase of rotational speed of runner, the load carrying capacity is found to increase with the speed up to a certain speed for specified parameters in table 5.1. This may be due to the assumption taken and at higher speed the viscous force of air plays an important role.



**Fig 5.4: Variation of Load Carrying Capacity with Speed of the Runner**

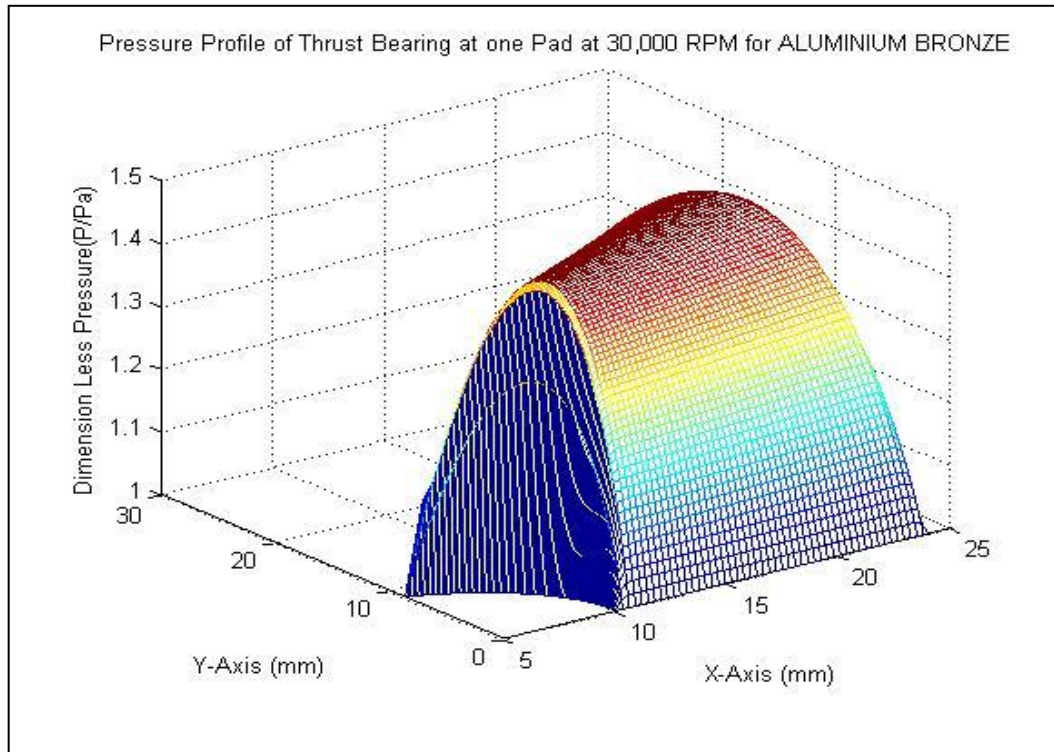
Figure 5.5 shows the variation of load carrying capacity with variation of thickness of bump foil from 100 microns to 800 microns. The load carrying capacity is found to be increased with the increase of thickness, however damping of the rotor may be affected at higher thickness of bump foils.



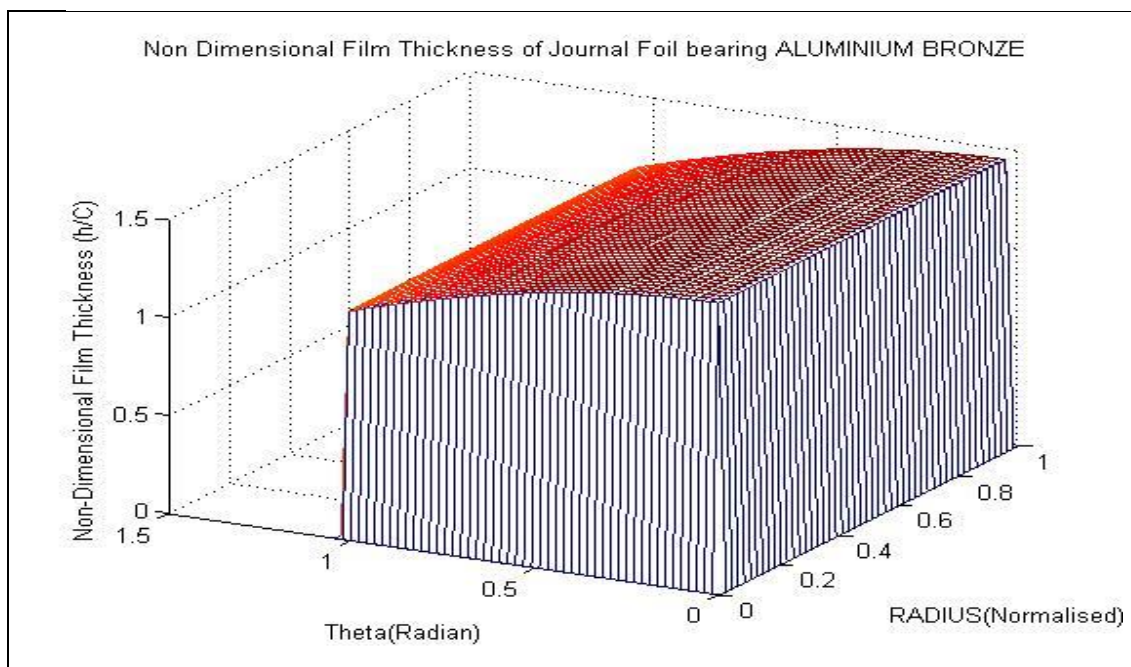
**Fig 5.5: Variation of Load Carrying Capacity with thickness of bump foil**

### 5.1 Analysis with Aluminium bronze

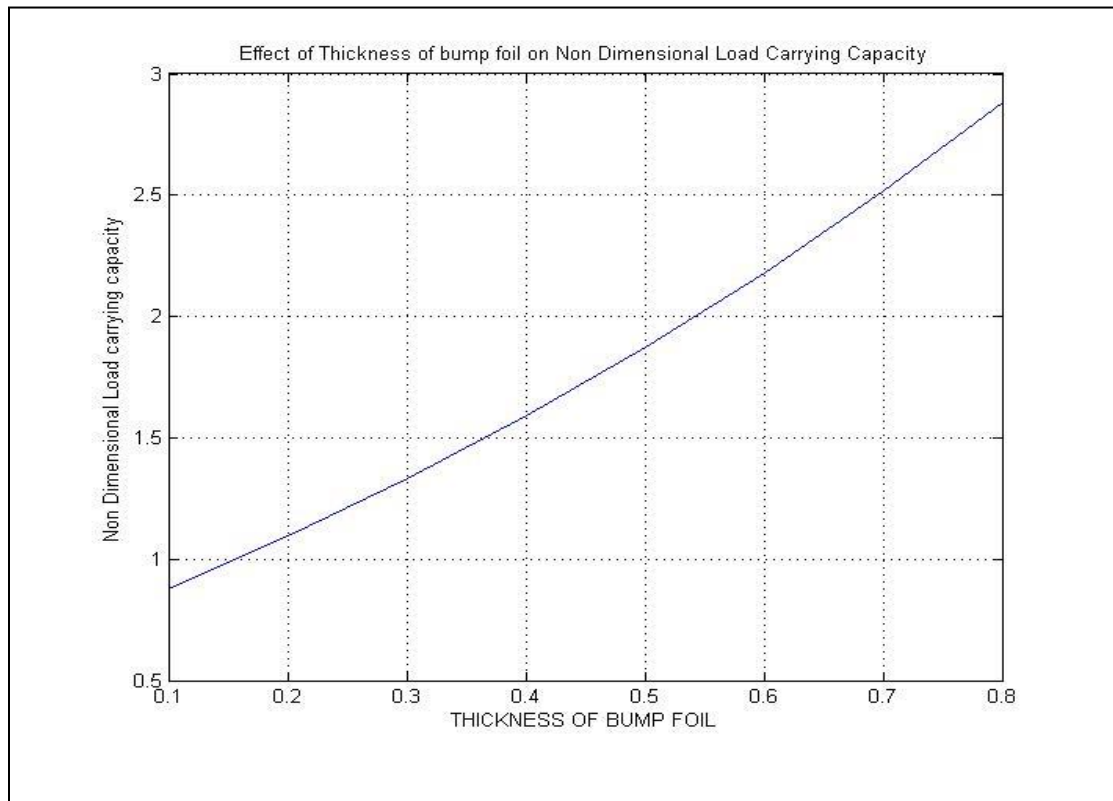
Similar analysis was done by changing only material of the bump foil. The elastic modulus of Aluminium bronze was taken as 120 GPa and poisson's ratio is 0.34. The result of analysis is shown from figure 5.5 to 5.10.



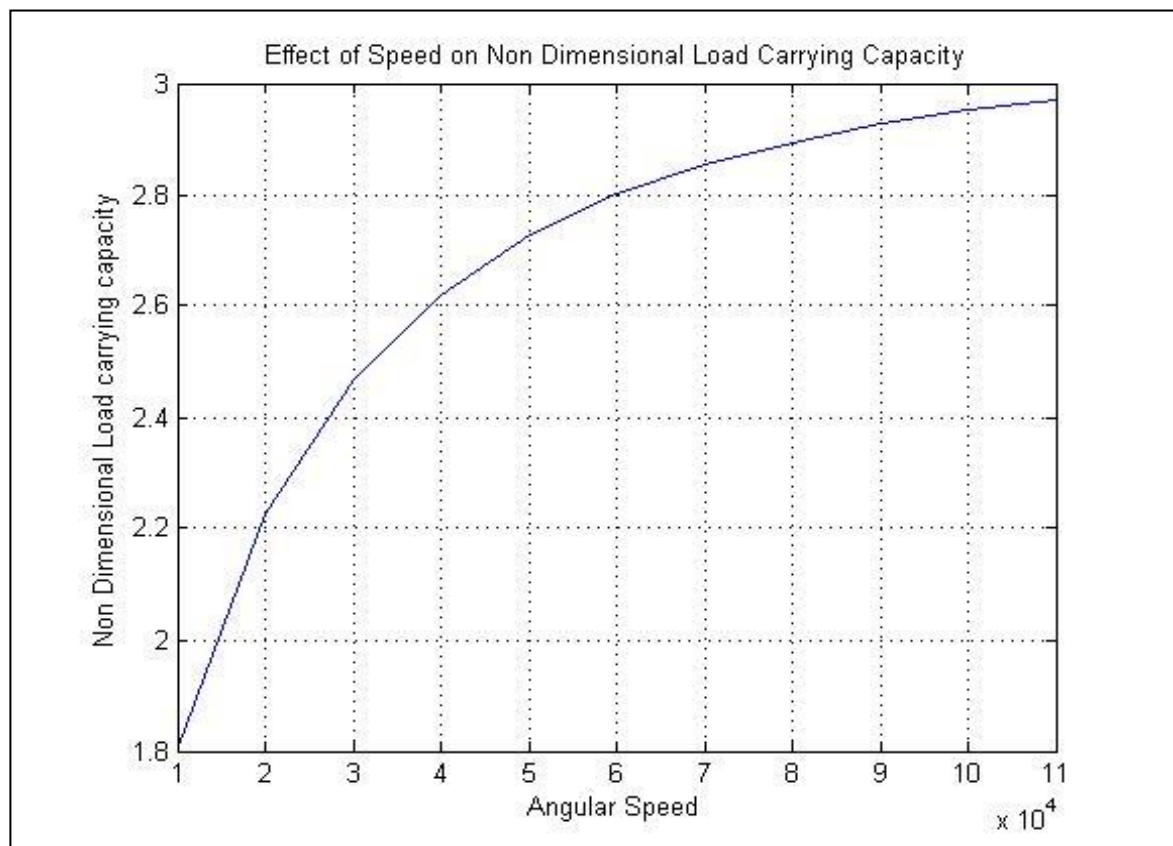
**Fig 5.6: Dimensionless pressure profile over one pad at 30,000 RPM**



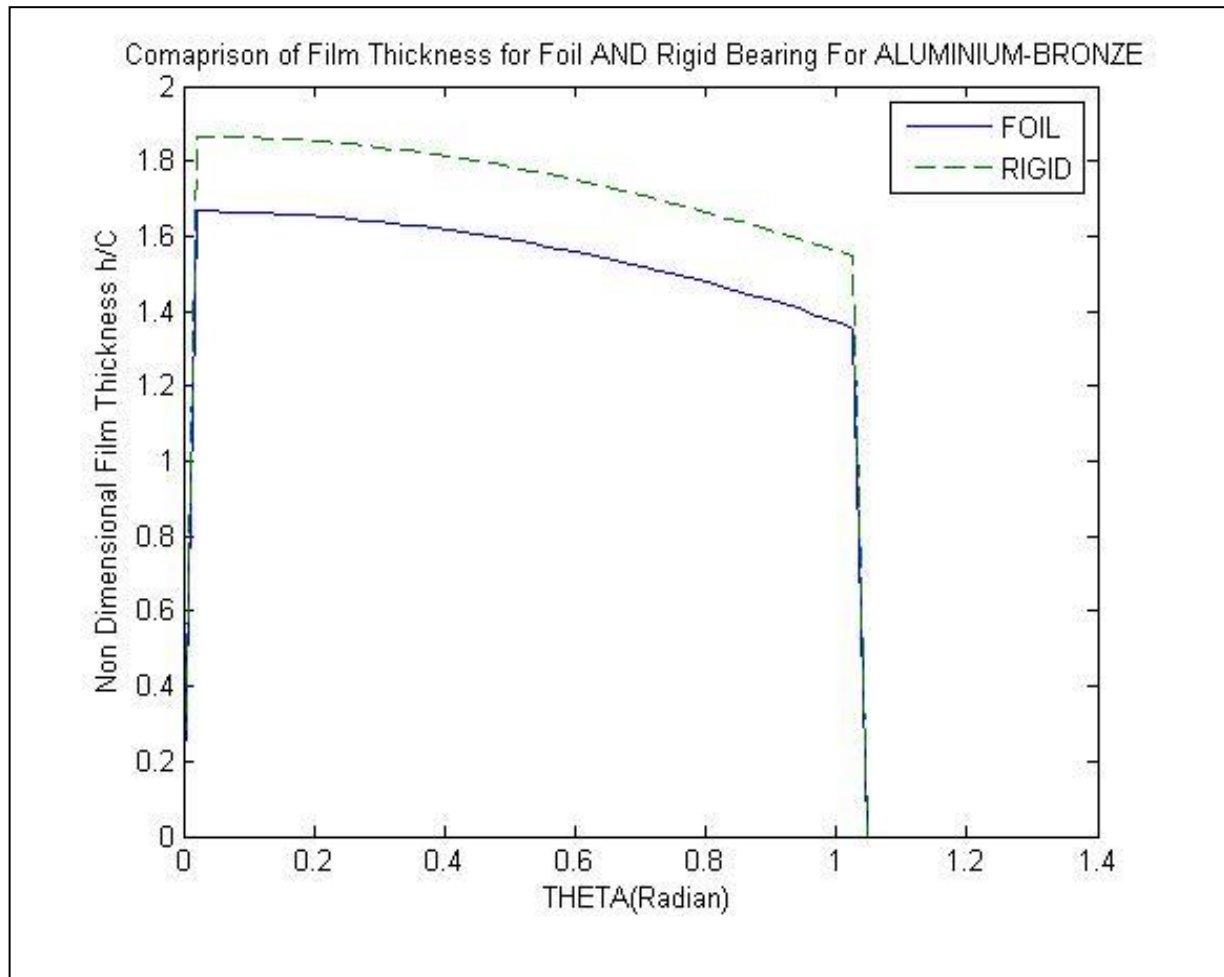
**Fig 5.7: Dimensionless Film thickness over one pad at 30,000 RPM**



**Fig 5.8: Variation of Load Carrying Capacity with thickness of bump foil**



**Fig 5.9: Variation of Load Carrying Capacity with Speed of the Runner.**



**Fig 5.10: Comparison of Film thickness between foil and rigid bearings over one pad at 30,000 RPM**

## **CHAPTER 6**

### **CONCLUSIONS**

The aerodynamic analysis of compliant foil thrust bearings helps to calculate the performance parameters. The compliant foil thrust bearings are function of fluid equation as well structural equating. So calculating actual performance parameters are difficult and consume lots of efforts. In the current project an attempt was made to find the performance parameters like pressure profile, film thickness load carrying capacity etc with several assumptions to simplify modified Reynolds equation. The analysis was done on two types of foils Inconel X 750 and Aluminium bronze. The Load carrying capacity of both types foils are found to satisfactory for use in small turbines and turbo generator. Due to protected technology of gas foil thrust bearings, very less no of work is found in open literatures. The result of current project may be help full to the researchers to work more on compliant foil thrust bearings.

### **FUTURE SCOPES**

- Experimental work can be carried out by designing the thrust bearings and testing it for results and comparing it with theoretical results.
- Stiffness and damping characteristics can be analytically found out which also influences load carrying capacity.
- Investigations can be carried out to find the optimum thickness of the foil and other geometry of bearings.
- Thermodynamic analysis of the foil bearings can be done to improve performance of bearing.

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13. **Robert J. Bruckner**(2012), “Performance of Simple Gas Foil Thrust Bearings in Air”,  
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14. **Robert Jack Bruckner**(2004), “Simulation and Modelling of Hydrodynamic, Thermal  
and Structural Behaviour of Foil Thrust Bearings”, Submitted for Ph.D. dissertation.  
Case Western Reserve University, School of Graduate Studies



## ROAD MAP

### Work done in 7<sup>th</sup> Semester

Schedule of work	Aug-12	Sep-12	Oct-12	Nov-12	Status
Literature survey					Completed
Derivation of Reynolds equation for Thrust bearings					Completed
Compressible Reynolds Equation in Two Dimensions					Completed
Finite difference analysis of 2D equations					Completed

### Work done in 8<sup>th</sup> Semester

Schedule of work	Jan-13	Feb-13	Mar-13	Apr-13	Status
Finite difference analysis of 2D equations					Completed
Writing program in MATLAB					Completed
Results and Discussions					Completed